

**2020 International Conference on Mathematical Neuroscience -  
Digital Edition (6th-7th of July 2020)**

***Dynamics of homogeneous and heterogeneous  
networks of Hodgkin-Huxley-type of models with  
bistability between bursting and silent states***

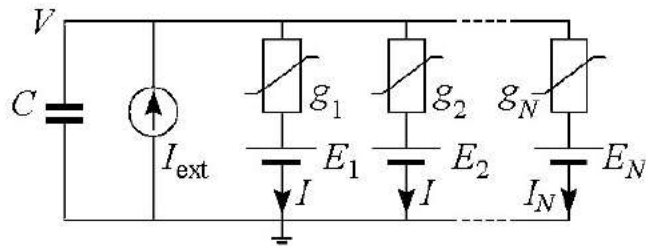
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**6 of July 2020**

# Hodgkin-Huxley formalism

The Hodgkin – Huxley model is a mathematical model that describes the generation and distribution of action potentials in neurons. Similar models were subsequently created for other electrically excited cells: cardiac myocytes, pancreatic beta cells.



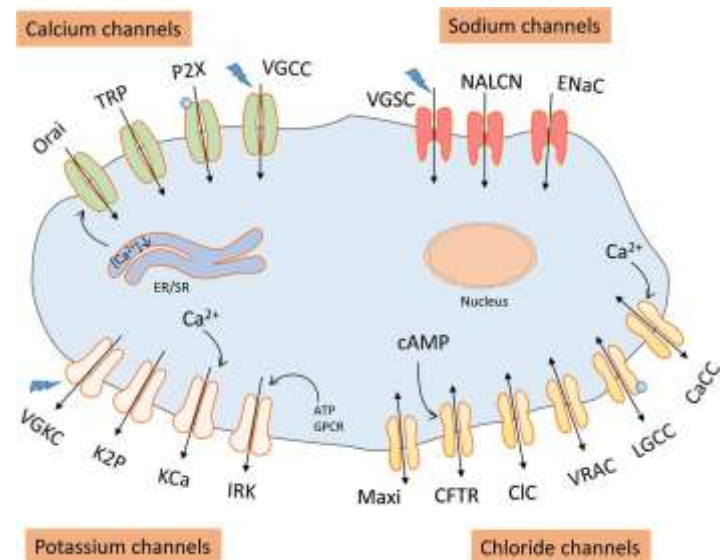
Equivalent scheme of neuron membrane

Hodgkin-Huxley model:

$$C \frac{dV}{dt} = I_{ext} - \sum_{j=1}^N I_j \quad I_j = g_j(V - E_j)$$

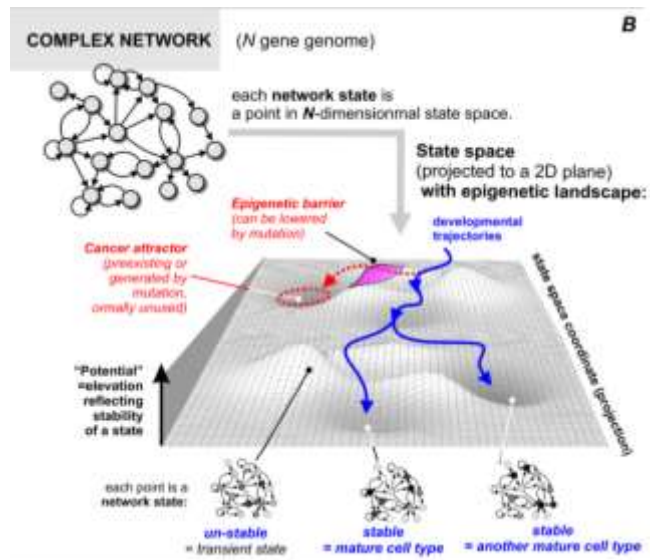
$$C \frac{dV}{dt} = I_{ext} - \bar{g}_K n^4 (V - V_K) - \bar{g}_{Na} m^3 h (V - V_{Na}) - \bar{g}_L (V - V_L),$$

$$\frac{dx}{dt} = \alpha_x(V)(1 - x) - \beta_x(V)x.$$



# Multistability is a ubiquitous Phenomenon

Multistability of oscillatory and silent regimes is a ubiquitous phenomenon exhibited by excitable systems such as neurons, cardiac cells, pancreatic beta-cells etc.



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## Cancer attractors: A systems view of tumors from a gene network dynamics and developmental perspective

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## Dynamics inside the cancer cell attractor reveal cell heterogeneity, limits of stability, and escape

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Pisarchik A.N., Feudel U., Control of multistability, *Physics Reports*, 2014, vol. 540, pp. 167-218.

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Loewenstein, Y., Mahon, S., Chadderton, P., Kitamura, K., Sompolinsky, H., Yarom, Y., & Ha'usser, M., Bistability of cerebellar Purkinje cells modulated by sensory stimulation, *Nature neuroscience*, 2005, vol. 8, pp. 202-211.

T. Malashchenko, A. Shilnikov, and G. Cymbalyuk, "Six types of multistability in a neuronal model based on slow calcium current," *PloS One* 518 6(7), e21782 (2011).

# Hodgkin-Huxley-type of model with bursting

Sherman model

$$\tau \dot{V} = -I_{Ca}(V) - I_K(V, n) - I_S(V, S),$$

$$\tau \dot{n} = \sigma(n_\infty(V) - n),$$

$$\tau_S \dot{S} = S_\infty(V) - S.$$

$$I_{Ca}(V) = g_{Ca} m_\infty(V)(V - V_{Ca})$$

$$I_K(V, n) = g_K n(V - V_K),$$

$$I_S(V, n) = g_S S(V - V_K),$$

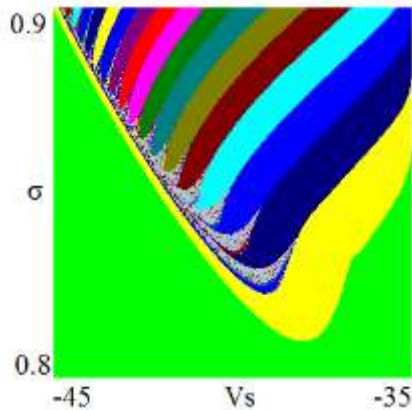


Chart of dynamical regimes

Table1. Parameters for original model

$\tau =$	0.02sec	$\tau_S =$	35sec	$\sigma =$	0.93
$g_{Ca} =$	3.6	$g_K =$	10.0	$g_S =$	4.0
$V_{Ca} =$	25.0mV	$V_K =$	-75mV		
$\theta_m =$	12.0mV	$\theta_n =$	5.6mV	$\theta_s =$	10.0mV
$V_m =$	-20.0mV	$V_n =$	-16.0mV	$V_s =$	-35mV

$$\omega_\infty(V) = [1 + \exp \frac{V_\omega - V}{\theta_\omega}]^{-1}, \quad \omega = m, n, S$$

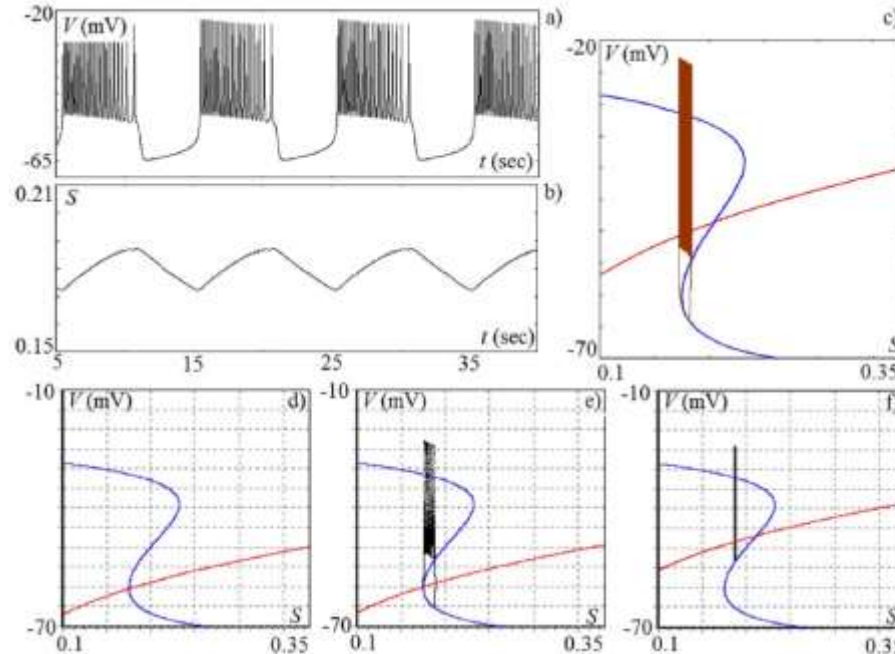


FIG. 1. Time series of the fast (a) and slow (b) variables; (c) fast (blue) and slow (red) manifolds together with a two-dimensional projection of phase portrait for the original Sherman model.<sup>18</sup> During the spiking phase,  $Ca^{2+}$ -ions flow into the cells and, during the silent phases,  $Ca^{2+}$ -ions are pumped out. The fast (spiking) dynamics is related to the flow of  $K^+$ -ions. (d)  $V_S = -45$  mV; (e)  $V_S = -44.7$  mV; and (f)  $V_S = -33.7$  mV.

# Bistability between silent and bursting states

Stankevich N.V., Mosekilde E. Coexistence between silent and bursting states in a biophysical Hodgkin-Huxley-type of model // CHAOS. 2017. Vol. 27. Issue 11.

Modified model with new  $K^+$  ion channel

$$\tau \dot{V} = -I_{Ca}(V) - I_K(V, n) - I_{K2}(V) - I_S(V, S),$$

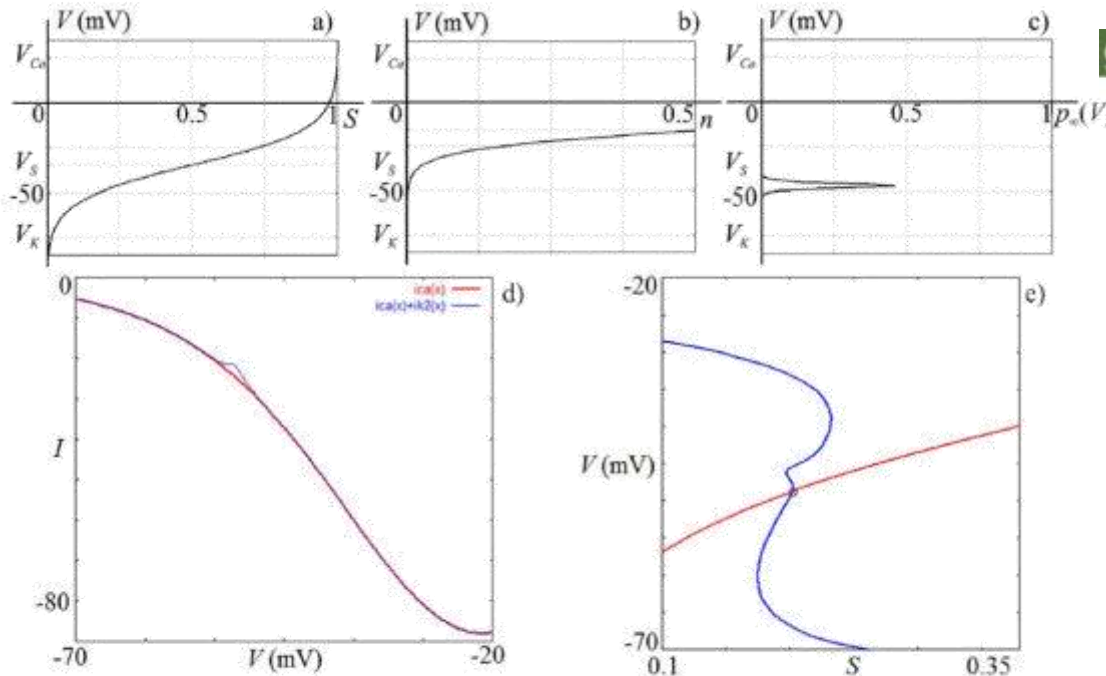
$$\tau \dot{n} = \sigma(n_\infty(V) - n),$$

$$\tau_S \dot{S} = S_\infty(V) - S.$$

EP (-49.084, 0.0027105, 0.19648)

$$I_{K2}(V) = g_{K2} p_\infty(V) (V - V_K)$$

$$p_\infty(V) = \left[ \exp \frac{V - V_p}{\theta_p} + \exp \frac{V_p - V}{\theta_p} \right]^{-1}$$



Review

## Neuronal Voltage-Gated Calcium Channels: Structure, Function, and Dysfunction

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Voltage-gated calcium channels are the primary mediators of calcium entry into neurons. There is great diversity of calcium channel subtypes across all animals, consistent with a variety of underlying signaling. This allows these channels to fulfill highly specialized particular subcellular tasks. While calcium channels are of critical importance in neurons, their dysfunction gives rise to a variety of neurodegenerative and psychiatric disorders. This Review discusses selected aspects of calcium channel structure, function, and dysfunction.

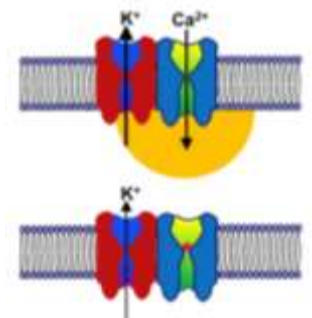
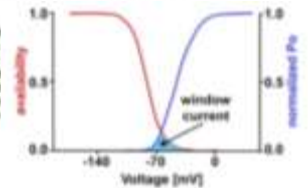
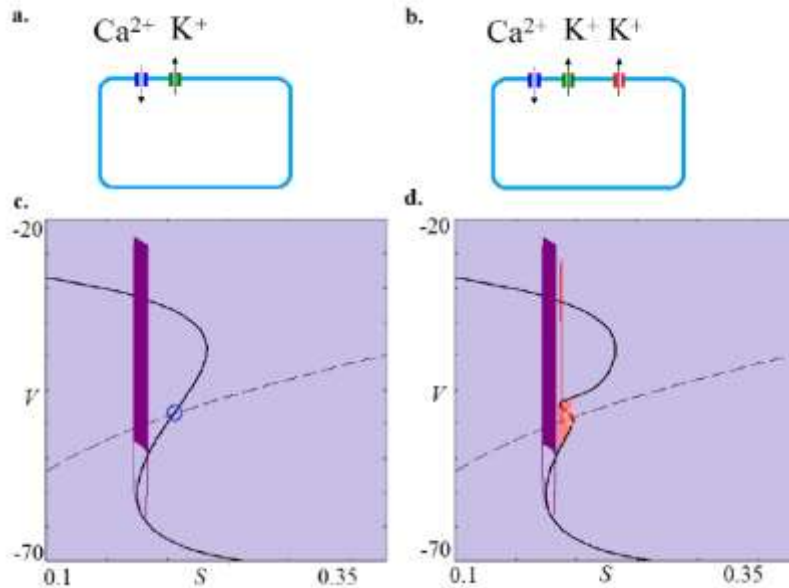


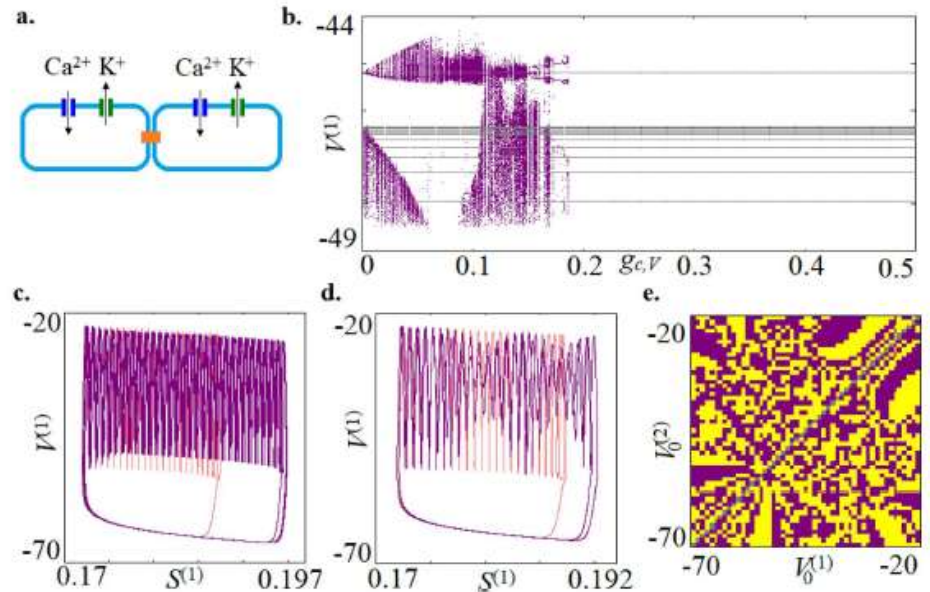
FIG. 2. Dependence of the membrane potential on the different ions: (a) calcium channel; (b) potassium channel; (c) probability function of new ion channel; (d) current of  $Ca^{2+}$  (2) (red) and sum of current  $Ca^{2+}$  (2) and current of new channel (6) (blue); (e) fast (blue) and slow (red) manifolds of the modified model (8). Supplementing parameters for the new ion channel:  $g_{K2}=0.14$ ,  $\theta_p=1$  mV, and  $V_p=-46$  mV.

# The simplest homogeneous networks:



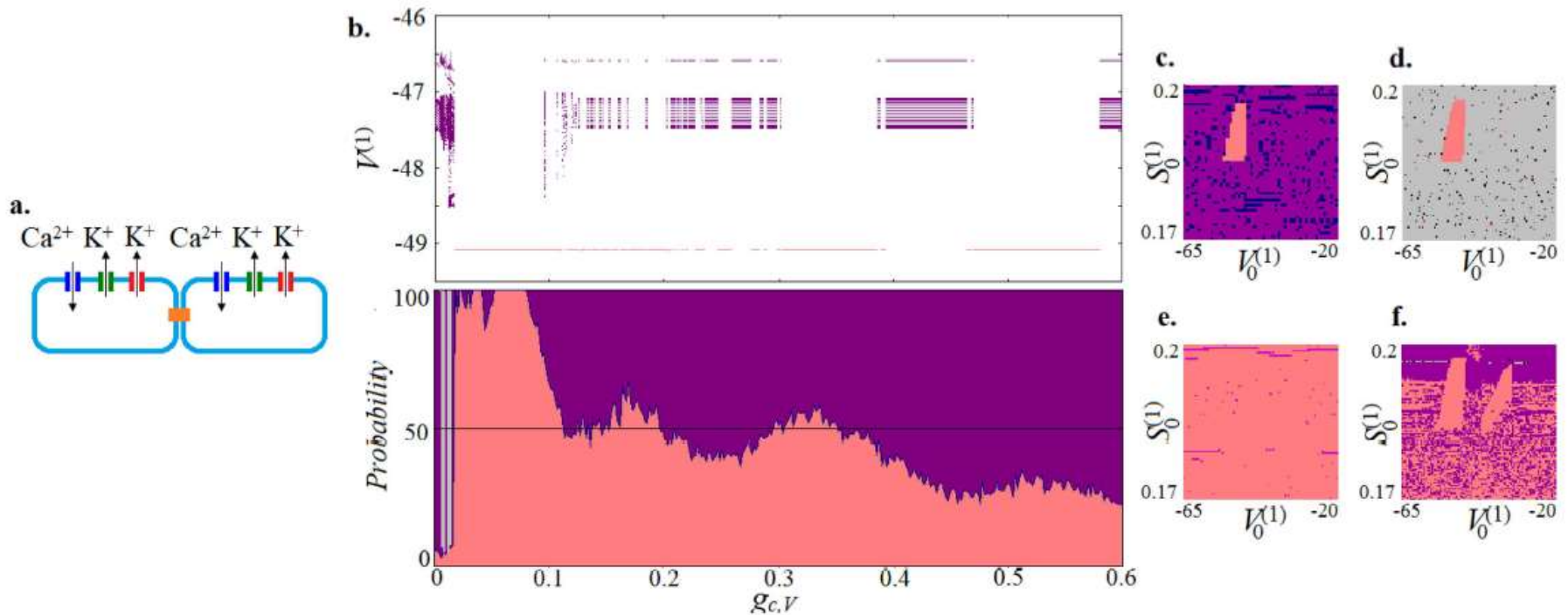
$$I_C(V^{(i)}) = \sum_{j \in \Gamma_i} g_{c,V}(V^{(i)} - V^{(j)}),$$

**Typical dynamical behavior for coupled identical oscillators: multistability and chaos for small coupling and synchronization for larger coupling**



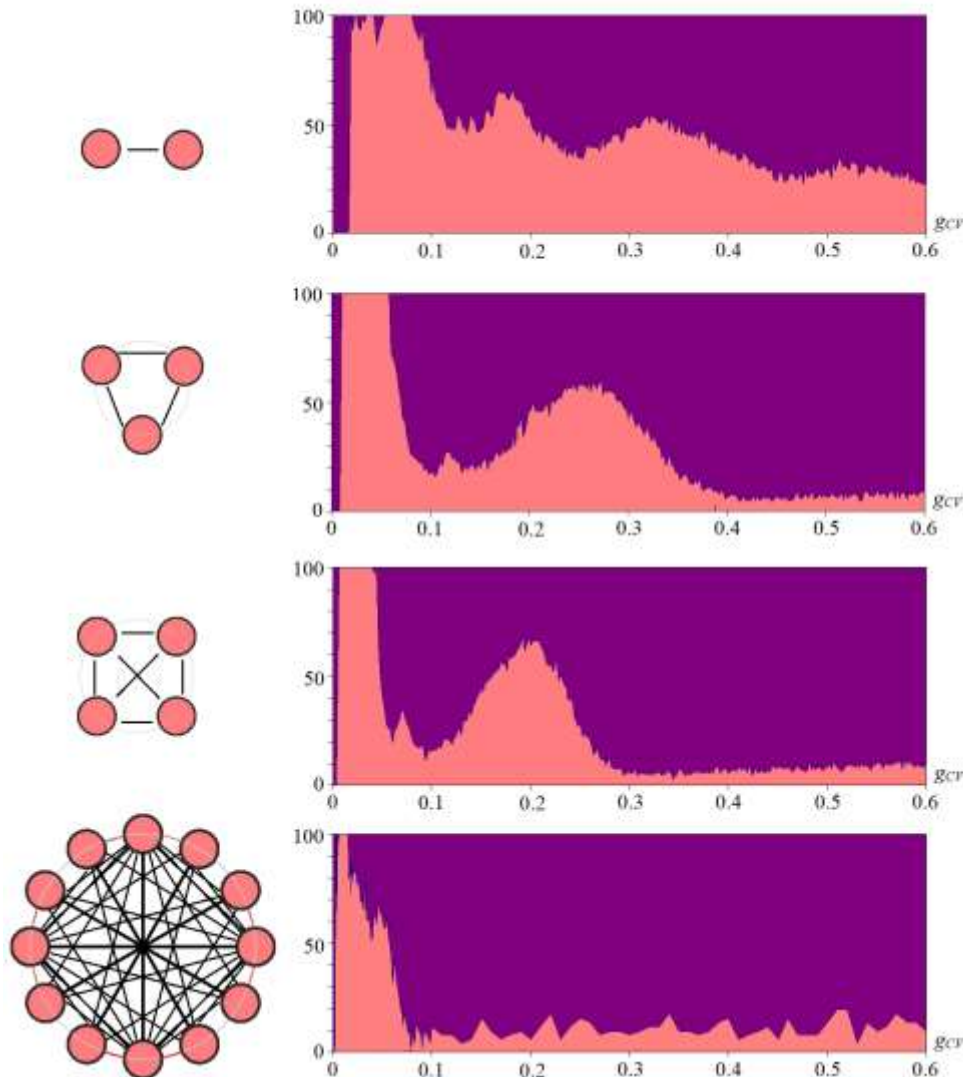
**FIG. 2.** Behavior of two coupled  $\beta$ -cells with bursting dynamics. (a) Schematic representation of the model. (b) Bifurcation trees estimated from Poincaré sections at the hypersurface  $n^{(1)} = 0.003$  when  $g_{c,V}$  is scanned in both directions (violet: forward, gray: backward).  $k^{(i)} = 0, i = 1, 2$  and other parameters as in Fig. 1. Exemplary phase portraits depicting coexistence between synchronized bursting and (c) periodic or (d) aperiodic attractors are also shown for  $g_{c,V} = 0.03$  and  $0.15$ , respectively. (e) Estimated basins of attraction of the coexisting bursting attractors for  $g_{c,V} = 0.15$  when starting from the following initial conditions:  $n_0^{(1)} = 0.003, S_0^{(1)} = 0.1889747, n_0^{(2)} = 0.0025039, S_0^{(2)} = 0.1889566$ .

# The simplest homogeneous network with bistability



**FIG. 3.** Behavior of two coupled  $\beta$ -cells with dysfunctional potassium channels. (a) Schematic representation of the two-cell system. (b) (Top) Bifurcation tree estimated from Poincaré sections at the hypersurface and  $n^{(1)} = 0.003$  constructed for fixed initial conditions:  $V_0^{(1)} = -50.784$ ,  $n_0^{(1)} = 0.0245$ ,  $S_0^{(1)} = 0.1774$ ,  $V_0^{(2)} = -49.084$ ,  $n_0^{(2)} = 0.027105$ ,  $S_0^{(2)} = 0.19648$ , as well as (bottom) probability of the appearance of coexisting attractors in dependence on the coupling strength  $g_{c,V}$ .  $k^{(i)} = 1$ ,  $i = 1, 2$  and other parameters as in Fig. 1. Exemplary basins of coexisting attractors for various strength of coupling: (c)  $g_{c,V} = 0.004$ ; (d)  $g_{c,V} = 0.0063$ ; (e)  $g_{c,V} = 0.02$ ; (f)  $g_{c,V} = 0.12$ . For (c)–(f), the initial conditions of the other variables were fixed in steady state  $n_0^{(1)} = 0.027105$ ,  $V_0^{(2)} = -49.084$ ,  $n_0^{(2)} = 0.027105$ ,  $S_0^{(2)} = 0.19648$ . Gray: aperiodic dynamics; violet: bursting dynamics; red: steady state.

# Larger homogeneous networks with bistability

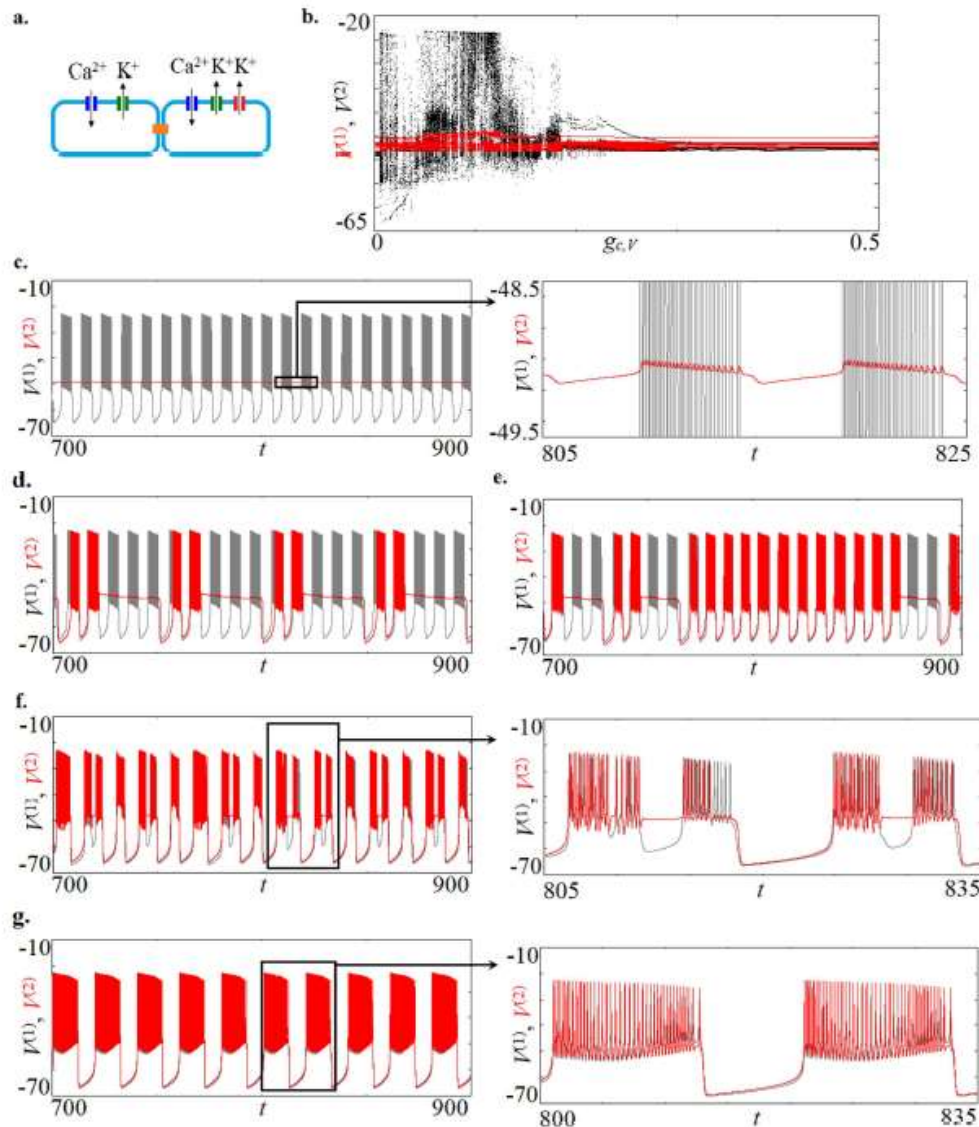


**Fig.** Behavior of network of coupled beta-cells with dysfunctional potassium channels. (a) Schematic representation of the network-cell system. (b) Probability of the appearance of coexisting attractors in dependence on the coupling strength  $g_c$ ,  $V_i$ ,  $k(i) = 1$ ,  $i = 1, 2$  and other parameters as in Fig. 1.

***Dominance of silent dynamics for certain level of coupling strength***



# The simplest heterogeneous networks with bistability

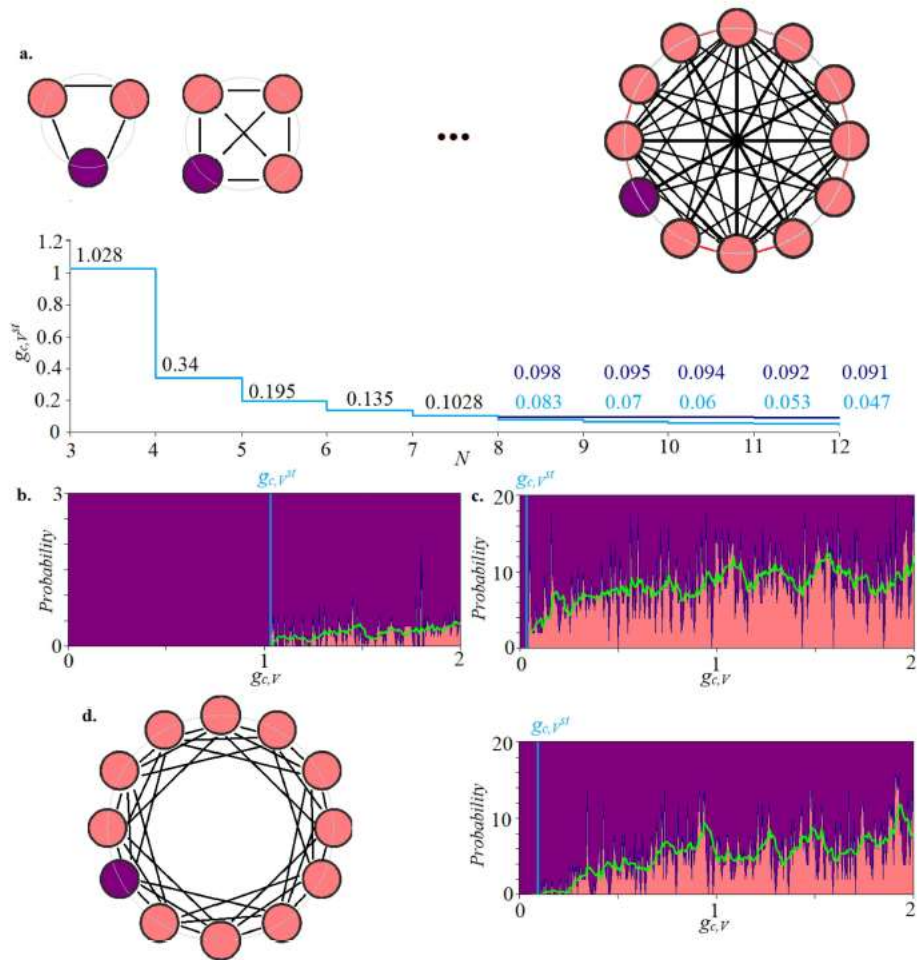


**For minimal network  
silent state is not  
stabilized.**

**Mix-mode dynamics is  
observed for small  
coupling strength.  
Synchronization for  
larger coupling.**

**FIG. 4.** Behavior of a minimal mixed population model. (a) Schematic representation of minimal mixed population of two coupled cells (only one cell has channel dysregulation). (b) Bifurcation tree estimated from Poincaré sections at the hypersurface  $n^{(1)} = 0.003$  when  $g_{c,v}$  is scanned in both directions.  $k^{(1)} = 0, k^{(2)} = 1$ , and other parameters as in Fig. 1. Exemplary time series of characteristic dynamical regimes for (c)  $g_{c,v} = 0.001$ , (d)  $g_{c,v} = 0.008$ , (e)  $g_{c,v} = 0.015$ , (f)  $g_{c,v} = 0.07$ , (g)  $g_{c,v} = 0.3$ .

# Larger heterogeneous networks with bistability



**FIG. 5.** Behavior of mixed  $\beta$ -cells population. (a) (Top) Schematic representation of globally coupled networks of increasing size, where only a single cell does not have a channel dysregulation (red color of node responds to  $k^{(i)} = 1$ , violet color to  $k^{(i)} = 0$ ) and (bottom) a corresponding plots of coupling strength thresholds for steady state stabilization in dependence on the size of the population ( $N$ ) for globally (light blue) and locally (dark blue) coupled cells. Schematic representation of locally coupled population in Fig. 5(d), left. Probability of the appearance of coexisting attractors in dependence on the coupling strength for the minimal population (b) of  $N = 3$  coupled cells, as well as for (c)  $N = 12$  globally or (d) locally coupled cells. Other parameters as in Fig. 1.  $g_{c,v}^{st}$  -  $g_c, v$ , threshold for which the steady state is stabilized.

**Stabilization of equilibrium is possible for larger networks, when number of oscillators with bistability more than  $N/2$ ,  $N$  is size of network. Effect for globally and locally coupled network is the same.**

## ***Conclusion***

***Homogenous networks with bistability demonstrates dominance of silent dynamics for certain level of coupling.***

***In heterogeneous networks stabilization of equilibrium is possible, when number of oscillators with bistability more than 50% in network.***

***Stabilized silent state in heterogeneous networks doesn't dominate, possibility to obtain such kind of behavior is less than 20%***

***Effect for globally and locally coupled networks is the same.***